

Thin Film Conductivity Measurements

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Reminder of Experimental Setup

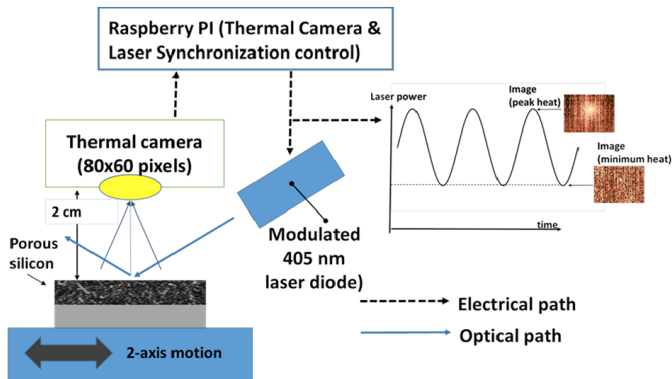


Figure 1: Optical system diagram for the thermography system developed as part of this work to investigate thermal characteristics of porous silicon. The physical housing of the system within a $20 \times 20 \times 13 \text{ cm}^3$ IP66 rated enclosure.

Approach Taken

- ▶ Understand fundamentals of actual situation
- ▶ Attack a simplified problem
 - ▶ 1-D version
 - ▶ No Heat loss to substrate or air
- ▶ Model heat loss in 1-D version

Physical Setup

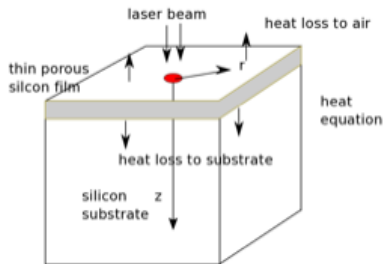


Figure 1: Simplified view of experiment

Fundamental equations

$$\frac{\partial T}{\partial t} = \kappa_f \nabla^2 T + \frac{Q_0}{h} \delta(r) e^{i\omega t} \quad \text{Film}$$

$$\frac{\partial T}{\partial t} = \kappa_s \nabla^2 T \quad \text{Substrate}$$

Complications

- ▶ Thin Layer
- ▶ Conductivity Ratio About 100 : 1
- ▶ Hard To Asses Heat Loss To Atmosphere

Aim

How best to determine thermal properties of film
(what ω to use

Wolf's arrangement

A simpler arrangement has been explored by Wolf.

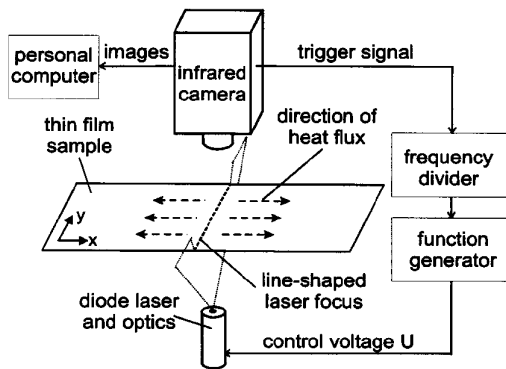


FIG. 1. Experimental setup: A sinusoidally modulated diode laser illuminates the backside of a thin film sample. The focus of the laser light is line shaped. The induced thermal wave is visualized by a lock-in procedure using an IR camera. The dashed arrows symbolize the direction of the heat flux.

Figure 2: Wolf cuts off the film and heats along a line: much harder experimentally, *much* easier analytically

Equations

$$\text{Heat Equation : } \rho_f C_f T_t(x, t) = k_f T_{xx}(x, t) \quad (1)$$

$$\text{Heat Input : } -k_f T_x(x, t) \Big|_{x=0} = q_0 e^{i\omega t} \quad (2)$$

We expect a steady state (or oscillatory) solution of the form

$$T(x, t) = e^{i\omega t} X(x) \quad (3)$$

Scaling

Set

$$\begin{aligned}T &= T_o + \Delta T T' \\ \omega t &= t' \\ x &= x_0 x'\end{aligned}\tag{4}$$

Substituting (4) into (1) we obtain

$$\rho_f C_f \Delta T \omega = \frac{k_f \Delta T}{x_0^2} T'_{x'x'}\tag{5}$$

This leads to

$$x_0^2 = \frac{K_f}{\omega} \Rightarrow x_0 = \sqrt{\frac{K_f}{\omega}}\tag{6}$$

with K_f the diffusivity given by

$$K_f = \frac{k_f}{\rho_f C_f}\tag{7}$$

The scaled result is

$$T'_{t'} = T'_{x'x'}\tag{8}$$

Scaling

Substituting (4) into (2) we obtain

$$-\frac{k_f \Delta T}{x_0} \frac{\partial T'}{\partial x'} = q_0 e^{it'} \quad (9)$$

This leads to

$$\Delta T = \frac{q_0}{\sqrt{k_f \rho_f C_f \omega}}$$

The result of scaling is

$$\left. \frac{\partial T'}{\partial x'} \right|_{x=0} = -e^{it'}$$

Observations

$$x_0 \propto \frac{1}{\sqrt{\omega}}$$

$$x_0 \propto \sqrt{K_f}$$

$$\Delta T \propto \frac{1}{\sqrt{\omega}}$$

$$\Delta T \propto q_0$$

The substitution of (3) into (8) gives

$$X_{xx} - iX = 0 \quad (10)$$

with

$$X(x) = Ae^{-\frac{1}{\sqrt{2}}(1+i)x} \quad (11)$$

with Constant of integration A given by

$$A = \frac{\sqrt{2}}{(1+i)} \quad (12)$$

Thus,

$$X(x) = \frac{\sqrt{2}}{(1+i)} e^{-\frac{1}{\sqrt{2}}(1+i)x} \quad (13)$$

Hence,

$$T(x, t) = e^{\frac{-x}{\sqrt{2}}} e^{i\left(t - \frac{x}{\sqrt{2}} - \frac{\pi}{4}\right)} \quad (14)$$

Results For No Heat Loss Case (*Insulated film*)

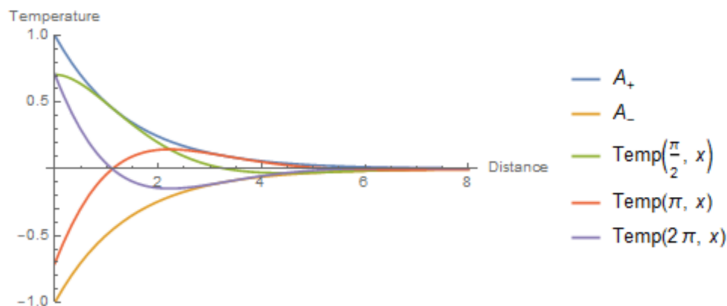


Figure 3: Temperature Profiles for different t

Heat Losses Case

$$\rho_f C_f T_t = k_f T_{xx} - \gamma(T - T_0) \quad (15)$$

$$-k_f T_x = q_0 e^{i\omega t} \quad (16)$$

with the same scaling we obtained

$$\begin{aligned} T'_{t'} &= T'_{x'x'} - \mu T' \\ -T'_{x'} &= e^{it} \end{aligned}$$

where μ is scaled heat loss parameter and it is

$$\mu = \frac{\gamma}{\rho_f C_f \omega}$$

Thus,

$$X(x) = \frac{1}{\sqrt{i + \mu}} e^{(\sqrt{i + \mu})x} \quad (17)$$

$$T(x, t) = e^{it} \frac{1}{\sqrt{i + \mu}} e^{(\sqrt{i + \mu})x} \quad (18)$$

$$\sqrt{\mu + i} = \sqrt{r_0} \left[\cos \left(\frac{\pi}{4} - \frac{\phi}{2} \right) + i \sin \left(\frac{\pi}{4} - \frac{\phi}{2} \right) \right] \quad (19)$$

where $\tan(\phi) = \mu$.

Hence,

$$X(x) = \frac{1}{\sqrt{\mu + i}} \left[e^{-\sqrt{r_0} \left[\cos \left(\frac{\pi}{4} - \frac{\phi}{2} \right) x \right]} e^{-i \sqrt{r_0} \left[\sin \left(\frac{\pi}{4} - \frac{\phi}{2} \right) x \right]} \right] \quad (20)$$

As $\mu \rightarrow 0$, we go back to the simple case.

(20) approximately gives

$$AMP(x) = \left[1 - \left(\frac{\phi^2}{2} \right) \right] e^{\frac{-1}{\sqrt{2}} \left(1 + \frac{\phi}{2} \right) x} \quad (21)$$

Results For Heat Loss Case (*Not insulated*)

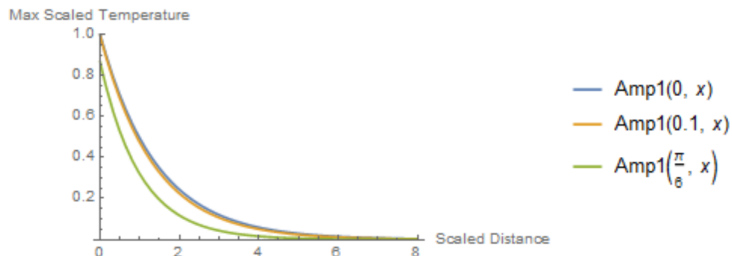


Figure 4: Temperature Profiles for different t

Conclusion

We presented the general problem and examined the situation in which heat was supplied along a line:

- ▶ In the special ideal case in which there was no heat loss, we found that waves travel away from the line source with decaying amplitude
- ▶ The maximum temperature change reached was

$$\frac{q_0}{\sqrt{k_f \rho_f C_f \omega}}$$

- ▶ These results can be used to determine K_f , k_f and $\rho_f C_f$
- ▶ We modelled effect of heat losses and found formulae that would be used to account for the heat losses
- ▶ If heat loss parameter is less than **0.3**, the effect is minor.