

Turbulent Fluid Fracture

8th January 2011

MISG 2011

Overview

- ① Problem Description
- ② Derivation
- ③ Analytical solution
- ④ Numerical solution
- ⑤ Comparisons
- ⑥ Conclusion
- ⑦ Q and A

Problem description

Descriptive Equation

$$\frac{\partial h}{\partial t} = -D^* \frac{\partial}{\partial x} \left[\left(-h^3 \frac{\partial h}{\partial x} \right)^{\frac{1}{m+2}} \right] \quad (1)$$

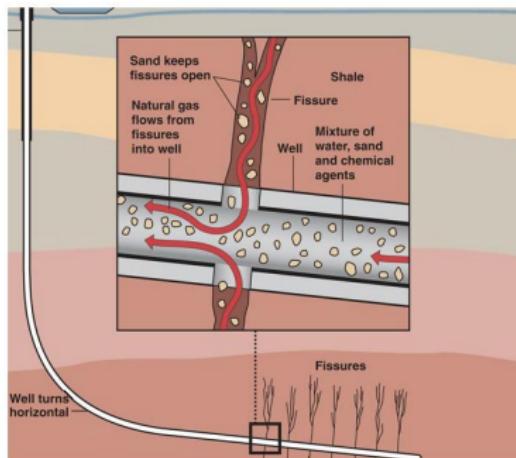


Figure: A

Derivation

$$\frac{2\tau_0}{\rho \bar{u}^2} = n \left[\frac{2\rho h \bar{u}}{\mu} \right]^m = n Re^m \quad \text{where } Re = \frac{2\rho h \bar{u}}{\mu} \quad (2)$$

$$\tau_0 = -h \frac{\partial p}{\partial x} \quad (3)$$

$$\frac{\partial h}{\partial t} + \frac{\partial}{\partial x}(h \bar{u}) = 0 \quad (4)$$

The use of equations (2) - (4), Assumption: $p(x, t) = \Lambda h(x, t)$

$$\frac{\partial h}{\partial t} = -D^* \frac{\partial}{\partial x} \left[\left(-h^3 \frac{\partial h}{\partial x} \right)^{\frac{1}{m+2}} \right] \quad (5)$$

$$\text{where } D^* = D \Lambda^{\frac{1}{m+2}} \text{ and } D = \left[\frac{2^{(1-m)} \mu^m \rho^{-(1+m)}}{n} \right]^{\frac{1}{m+2}}$$

Analytical solution

Reduction: PDE to ODE

- Scaling Transformations

$$\bar{t} = \lambda^a t, \quad \bar{x} = \lambda^b x, \quad \bar{h} = \lambda^c h$$

- Invariance of the PDE

$$c = \frac{m+2}{m-2}a - \frac{m+3}{m-2}b$$

This results in Quasi linear equation

$$cf = bx \frac{\partial f}{\partial x} + at \frac{\partial f}{\partial t} \tag{6}$$

- General Solution to the Quasi linear equation

$$h = f(x, t) = t^{\frac{c}{a}} F(\xi) \quad (7)$$

- Transformation $\xi = \frac{x}{t^\alpha}$ where $\alpha = \frac{b}{a}$.
- ODE

$$D^* \frac{d}{d\xi} [(-F^3 \frac{dF}{d\xi})^{\frac{1}{m+2}}] - \alpha \frac{d}{d\xi} (\xi F) + \frac{m+2-5\alpha}{m-2} F = 0 \quad (8)$$

Simplification of ODE

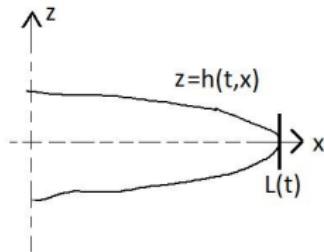


Figure: B

- moving boundary condition $L(t)$

First particular solution

- Elimination of the third term in eqn.:

Let $\alpha = \frac{1}{5}(m + 2)$

- Boundary condition

$$H(1) = 0$$

- Integration of eqn. w.r.t. η ,

Assumption: $(-H^3(1)H'(1))^{\frac{1}{m+2}}$

$$-H^3 \frac{dH}{d\eta} = (\alpha \eta H)^{m+2} \quad (10)$$

- Integration of eqn. q w.r.t. η , Condition: $m < 2$

$$H = [\alpha^{m+2} \frac{2-m}{m+3} (1 - \eta^{m+3})]^{\frac{1}{2-m}} \quad (11)$$

Comparison of Laminar and Turbulence cases:

Laminar	Turbulence
$m = -1$	$m = -\frac{1}{4}$
$\alpha = \frac{1}{5}$	$\alpha = \frac{7}{4}$
$L(t) = L_0 t^{\frac{1}{5}}$	$L(t) = L_0 t^{\frac{7}{4}}$

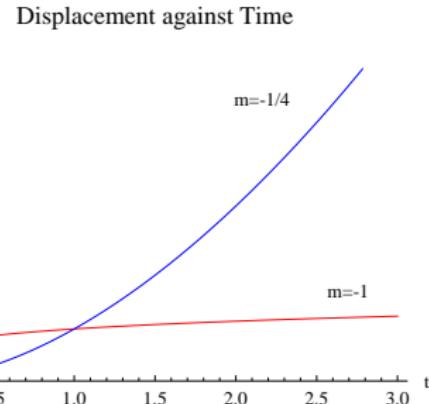
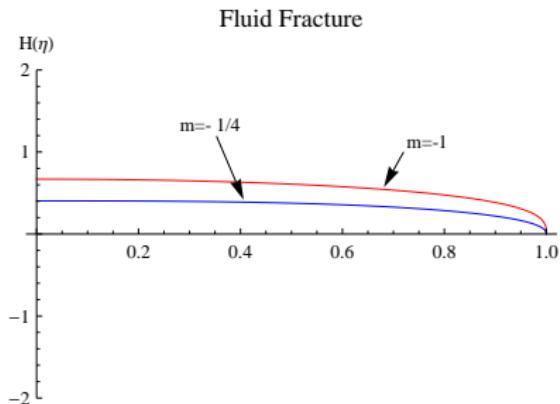


Figure: C

Second Solution

Ansatz:

$$H(\eta) = A(1 - \eta)^p \quad (12)$$

$$A = \left[\frac{(2-m)^{\frac{1}{m+2}}(m+2)}{m+4} \right]^{\frac{m+2}{2-m}},$$
$$p = \frac{1}{2-m}, \quad \alpha = 1.$$

Hence

$$H(\eta) = \left[\frac{(2-m)^{\frac{1}{m+2}}(m+2)}{m+4} \right]^{\frac{m+2}{2-m}} (1 - \eta)^{\frac{1}{2-m}} \quad (13)$$

Comparison of Laminar and Turbulence cases:

Laminar	Turbulence
$m = -1$	$m = -\frac{1}{4}$
$\alpha = 1$	$\alpha = 1$
$L(t) = L_0 t$	$L(t) = L_0 t$

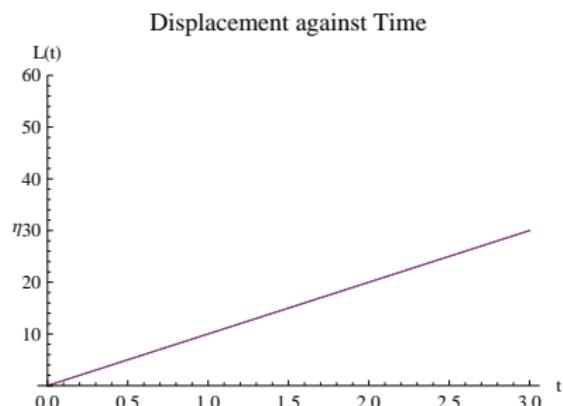
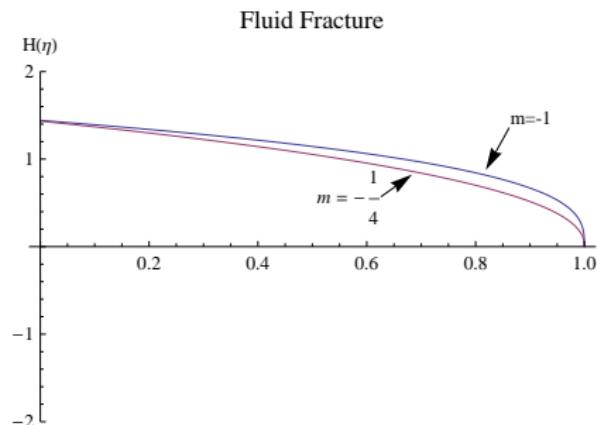


Figure: D

Conclusion

- Turbulence case appears to be fracturing faster than the laminar case.
 - Value of m can give us more information about the comparison of the two cases.
 - Phase plots could help in analysis of the stability of the solutions obtained.

Q and A